

Examples

Lawler Ch. 3

Karlin + Taylor Ch. 4

- 1) Birth-death processes where A is tridiagonal

$$A_{i, i+1} = \lambda_i = \text{aggregate birth rate}$$

$$A_{i, i-1} = \mu_i = \text{aggregate death rate}$$

$$A_{i, i} = -\mu_i - \lambda_i \quad (\text{row sums} = 0)$$

$$A_{i, j} = 0 \quad \text{for } |i - j| \geq 2$$

- a) Poisson process (# counts/arrivals)

$$\lambda_i = \lambda, \quad \mu_i = 0$$

- b) Queueing theory with k servers

$$\lambda_i = \lambda \quad (\text{arrival rate})$$

$$\mu_i = i\mu$$

$$\mu_i = k\mu \quad \text{for } i \geq k$$

where $\mu = \frac{1}{T_{\text{service}}}$ → average time for service

- c) Population models

$$\lambda_i = i\lambda \quad \text{or} \quad \lambda_i = i^2\lambda \quad \text{or} \quad \lambda_i = \lambda_i (k-i)$$

$$\mu_i = i\mu \quad \text{or} \quad \mu_i = i^2\lambda \quad \text{or} \quad \lambda_i = i\lambda + r$$

↑
immigrants

Annoying technicality: continuous-time MC is not necessarily well-defined by the matrix A if state space is infinite.

- no problem if A is bounded

- have to possibly worry if A has unbounded entries (K+T p. 135)

- may be problems if $\mu_i \propto i^2$ and $\lambda_i \propto i^2$
and $\frac{\lambda_i}{\mu_i} \rightarrow 0$ as $i \rightarrow \infty$

2) Finite-state processes w/o linear ordering

- transitions between energy levels in QM

- machine switching between modes of operation

- biomolecule switch between conformational states

How do we compute statistical properties of the continuous-time MC from its infinitesimal generator A ?

First we'll look at $P_{ij}(t) = \text{Prob}(X(t+t')=j | X(t')=i)$

Recall $A = \left. \frac{dP}{dt} \right|_{t=0^+}$

If A exists, then

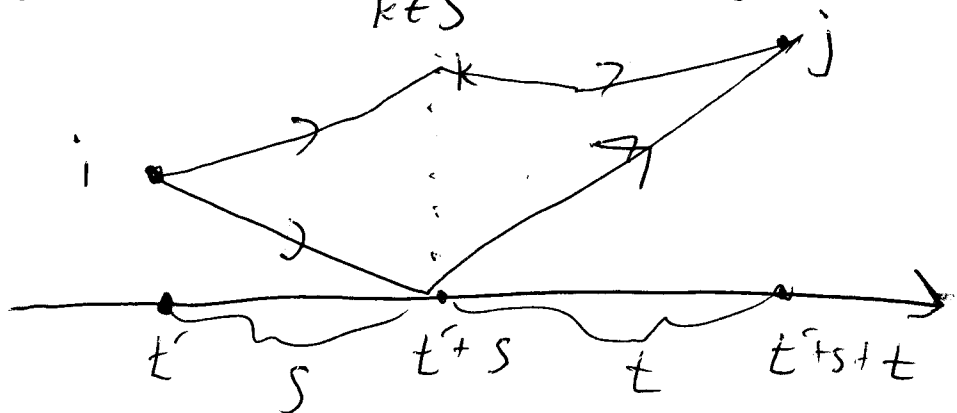
$$P_{ij}(\Delta t) = \delta_{ij} + \Delta t A_{ij} + o_{ij}(\Delta t)$$

↑
higher order terms
 $\frac{o_{ij}(\Delta t)}{\Delta t} \rightarrow 0$ as $\Delta t \rightarrow 0$

What about $P_{ij}(t)$ when t is not small?

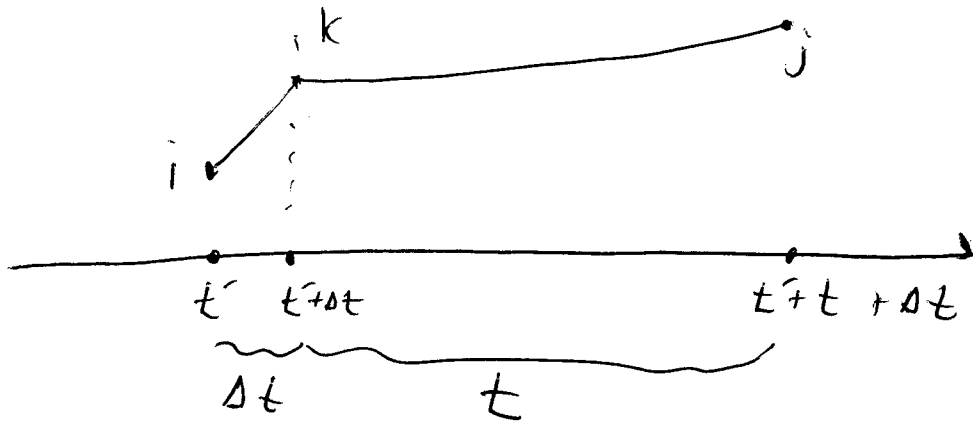
Chapman-Kolmogorov eqn in continuous-time:

$$P_{ij}(t+s) = \sum_{k \in S} P_{ik}(s) P_{kj}(t)$$



We'll use Chapman-Kolmogorov to recursively extend the time interval over which we know $P_{ij}(t)$.

$$P_{ij}(t + \Delta t) = \sum_{k \in S} P_{ik}(\Delta t) P_{kj}(t)$$



$$= \sum_{k \in S} (\delta_{ik} + \Delta t A_{ik} + o_{ik}(\Delta t)) P_{kj}(t)$$

$$P_{ij}(t + \Delta t) = P_{ij}(t) + \Delta t \sum_{k \in S} A_{ik} P_{kj}(t) + \sum_{k \in S} o_{ik}(\Delta t) P_{kj}(t)$$

$$\frac{P_{ij}(t + \Delta t) - P_{ij}(t)}{\Delta t} = \sum_{k \in S} A_{ik} P_{kj}(t) + \sum_{k \in S} \frac{o_{ik}(\Delta t)}{\Delta t} P_{kj}(t)$$

Formal $\Delta t \rightarrow 0$ limit:

$$\frac{d P_{ij}(t)}{dt} = \sum_{k \in S} A_{ik} P_{kj}(t)$$

$$\frac{d P}{dt} = A P$$

$$P(t=0) = I$$

Backward
Kolmogorov
equation

Rigorous handling of the $\Delta t \rightarrow 0$ limit:

- use dominated convergence theorem

with $P_{kj}(t) \leq 1$ and $\sum_{k \in S} a_{ik}(t) = \bar{a}_i(t) < \infty$

- Karlin + Taylor Vol. I + II

One could also repeat the argument starting from

$$P_{ij}(t) = \sum_{k \in S} P_{ik}(t) P_{kj}(\Delta t)$$

One gets from the formal calculation

$$\frac{d P_{ij}(t)}{dt} = \sum_{k \in S} P_{ik}(t) A_{kj}$$

$$\frac{d P}{dt} = P A$$

$$P(t=0) = I$$

Forward
Kolmogorov
Equation

The forward equation does not always make rigorous sense!

The solution to these equations (when they make sense) can be formally expressed as

$$P(t) = e^{At}$$

Makes perfect sense for finite state space

Several papers in SIAM Review

"Nine keen ways to compute Matrix
Exponentials"

Clive Maler and...

② Forward vs. Backward Equation

A) Consider first $\phi_j(t) = \text{Prob}(X(t) = j)$

$$\phi_j(t) = \text{Prob}(X(t) = j) = \sum_{k \in S} \text{Prob}(X(t) = j \text{ and } X(0) = k)$$

$$= \sum_{k \in S} \text{Prob}(X(t) = j | X(0) = k) \text{Prob}(X(0) = k)$$

$$\phi_j(t) = \sum_{k \in S} p_{kj}(t) \phi_k(0)$$

$$\vec{\phi}(t) = \vec{\phi}(0) \cdot P(t)$$

↑
vector

↑
matrix

Can avoid solving eqn for large matrix P by noting that $\vec{\phi}$ itself satisfies Kolmogorov forward eqn:

$$\boxed{\frac{d\vec{\phi}}{dt} = \vec{\phi} \cdot A} \quad (= \underbrace{\vec{\phi}(0) \cdot P(t)} \cdot A)$$

but not Kolmogorov backward eqn.

What is the Kolmogorov backward eqn good for?

Expectations starting from given state.

$$\text{Let } u_i(t) = \langle f(X(t)) | X(0) = i \rangle$$

$$= \sum_{j \in S} f(j) \text{Prob}(X(t) = j | X(0) = i)$$

$$= \sum_{j \in S} f(j) P_{ij}(t)$$

$$\vec{u}(t) = P(t) \cdot \vec{f}$$

This satisfies Kolmogorov backward equation:

$$\begin{aligned} \frac{d\vec{u}}{dt} &= A \vec{u} \\ \vec{u}(t=0) &= \vec{f} \end{aligned}$$

but not the forward eqn.

Forward equation describes how probabilities at a future time evolve.

Backward equation describes how expected "payouts" evolve given an initial state.

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Expectations starting from given state.

$$\begin{aligned} \text{Let } u_i(t) &= \langle f(X(t)) | X(0) = i \rangle \\ &= \sum_{j \in S} f(j) \text{Prob}(X(t) = j | X(0) = i) \\ &= \sum_{j \in S} f(j) P_{ij}(t) \\ \vec{u}(t) &= P(t) \cdot \vec{f} \end{aligned}$$

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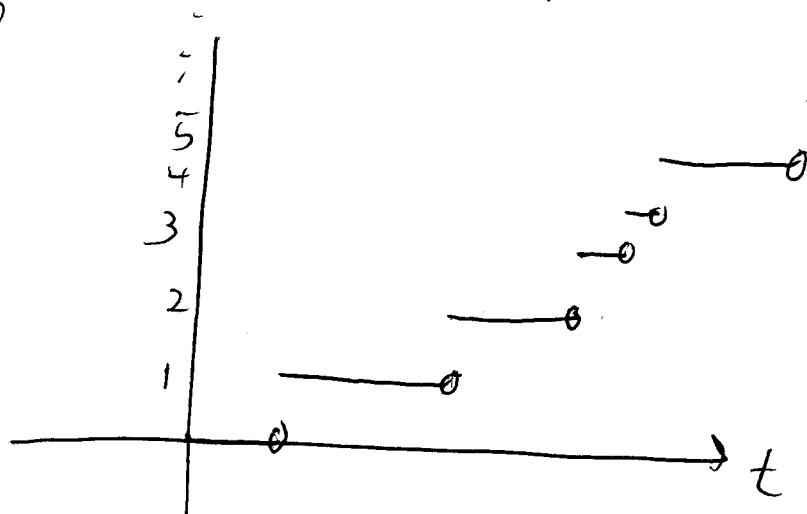
Poisson process

Poisson counting process $X(t)$ is a birth-death process with $X(0) = 0$, $\lambda_i = \lambda$, $\mu_i = 0$.

Infinitesimal generator

$$A = \begin{pmatrix} -\lambda & \lambda & & & \\ 0 & -\lambda & \lambda & & \\ & 0 & -\lambda & \lambda & \\ & & 0 & -\lambda & \lambda \\ & & & \ddots & \ddots \\ & & & & \ddots & \ddots \end{pmatrix}$$

What is the probability distribution for the amount of time $X(t)$ spends in a given state?



Consider $\phi_0(t) = \text{Prob}(X(t) = 0)$

Initialize: $\phi_0(0) = 1$.

$$\frac{d\vec{\phi}}{dt} = \vec{\phi} \cdot A \quad \text{Kolmogorov forward eqn}$$

0th component: $\frac{d\phi_0}{dt} = -\lambda \phi_0$

$$\phi_0(0) = 1$$
$$\phi_0(t) = e^{-\lambda t}$$

Let T be the random time at which $X(t)$ leaves state 0:

$$T = \inf_{t \geq 0} \{X(t) \neq 0\}$$

$$\begin{aligned} \text{Prob}(T > t) &= \phi_0(t) = \text{Prob}(X(t) = 0) \\ &= e^{-\lambda t} \end{aligned}$$

So if we write $\text{Prob}(T \in B) = \int_B p_T(t) dt$

then $p_T(t) = \lambda e^{-\lambda t}$

The transition time in a Poisson process has exponential distribution,
- same for any state

There is a deeper reason for why the time ~~to~~ to transition, T , has exponential distribution,

- Markov property that determines it,

$$\text{Prob}(T > t \mid T > s) = \text{Prob}(T > t - s)$$

(for $t > s$,
because given that $T > s$, there is
no memory about what happened before
time s ,)

$$\text{Prob}(X(t) = 0 \mid X(s) = 0, X(0) = 0)$$

$$\text{Prob}(X(t) = 0 \mid X(s) = 0)$$

By defn of cond. prob:

$$\frac{\text{Prob}(T > t)}{\text{Prob}(T > s)} = \text{Prob}(T > t - s)$$

$$\text{Let } G(t) = \text{Prob}(T > t)$$

$$*) \frac{G(t)}{G(s)} = G(t - s)$$

Clearly one solution is $G(t) = e^{-\lambda t}$ for some $\lambda > 0$.
This is the only class of solutions.

If $G(t)$ is differentiable, then

it differentiates $x)$ with respect to s ,
and write $u = t - s$:

$$G(t) = G(t-s) G(s)$$

$$0 = -G'(u) G(s) + G(u) G'(s)$$

$$\frac{G'(s)}{G(s)} = \frac{G'(u)}{G(u)} \quad \text{for all } s, u \geq 0.$$

$$\therefore \frac{G'(s)}{G(s)} = \text{constant}$$

There are also no nonsmooth solutions
Karlin + Taylor Th. 4.2.2,